

59<sup>th</sup> Rencontres de Moriond: QCD & High Energy Interactions, La Thuile, April 2025

### Theory Summary 2025

Peter Skands (Monash University)



EW/BSM  $Q \gtrsim M_{\rm EW}$ 

Fixed-Order pQCD Single-Scale  $Q \gg \Lambda$ 

Resummed pQCD Multi-Scale

 $Q_i \gg Q_j \gg \Lambda$ 

Hadronization (&UE) Powers, Strings, QGP  $Q \sim [\Lambda, 5\Lambda]$ 

> Hadrons  $Q \lesssim \Lambda$

> > $g_{\mu} - 2$

H, Z, X N latching Zt**, C** attice QED



### **Overview - 38 Theory Talks**

Warning!

Omissions, biases, and misrepresentations may be encountered beyond this point



## Higgs FLV & CPV



$$\begin{split} \mathcal{D}_{\widetilde{G}} &= f^{ABC} \widetilde{G}_{\mu}^{A\nu} \mathcal{G}_{\mu}^{B\nu} \mathcal{G}_{\rho}^{C\mu}, \qquad \mathcal{O}_{\widetilde{W}} = \epsilon^{IJK} \widetilde{W}_{\mu}^{I\nu} \mathcal{W} \qquad \overline{I}_{\mu}^{I} \mathcal{V} \qquad \mathcal{O}_{V}^{I} \mathcal{V}_{\mu}^{I} \mathcal{V}_{\mu}^{I}$$

 $\hat{O}_{tG} = g_{s} (\bar{Q}\sigma^{\mu\nu}T_{A}t)\tilde{\varphi}G^{A}_{\mu\nu}, \qquad \hat{O}_{tW} = i(\bar{Q}\sigma^{\mu\nu}\tau_{I}t)\tilde{\varphi}W^{\mu\nu}_{\mu\nu}, \qquad \hat{O}_{tW} = i(\bar{Q}\sigma^{\mu\nu}\tau_{I}t)\tilde{\varphi}W^{\mu$ 





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 $= \frac{y_t g_s \overline{t}_L \mathcal{S}_{Hj}^{\mu\nu} T^a G^a_{rata} (p)}{H_j (PT, for all of f$ 

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Moriond QCD 2025

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 $p_{T,H}$ 

## Drell-Yan

### Zhan: Intrinsic DY pT

FOPT divergent for  $p_T \rightarrow 0$ , calls for resummation

TMD PDF framework allows to resum logs of ( $p_T/M_Z$ )

Use MC "parton branching method" to evolve TMD PDFs



Starting point: assume Gaussian intrinsic component  $k_T/q_s$ : only affects low-p<sub>T</sub> part of spectrum

→ Extract  $q_s$  from ratio between low and high p<sub>T</sub> Low  $\sqrt{s} \lesssim 10 \,\text{GeV}: q_s \sim 1 \,\text{GeV}$ Higher  $\sqrt{s}: q_s$  increases

#### Vladimirov & Zhan



 $pp \rightarrow \gamma\gamma + jet \& pp \rightarrow pp$ 



03/04/20

#### Marcoli & Praszalowicz





## N3LO

#### **Pelloni:** 4-loop gg splitting functions



#### Hekhorn: aN3LO PDFs & Higgs Cross Sections



MSHTxNNPDFnnlo MSHTxNNPDFan3lo NNPDFan3lo MSHTan3lo PDF4LHC21 MSHTxNNPDFnnlo(qed) MSHTxNNPDFan3lo(qed) NNPDFan3lo(qed) MSHTan3lo(qed)

► NB: N<sup>3</sup>LO cross sections are a long term project ...



### Perturbative Uncertainties

Gocke (Mon): main uncertainties in top are ISR + FSR + "recoil" uncertainties (even with MiNNLOPS), JES (sens. to hadrochemistry), UE, and b-tagging Grohsjean (Mon): main uncertainties for toponium  $\eta_t$ : bb4l (offsh, tt & tW int), FSR modelling

**Poncelet (Thu):** missing higher orders in fixed-order and resummed predictions (MHOU)

Generic perturbative expansion:  $f(\alpha) = \alpha_s^n N_c^m \, \mathrm{d}\bar{\sigma}^{(0)} \left[ 1 + \alpha_s N_c \left( \frac{\mathrm{d}\bar{\sigma}^{(1)}}{\mathrm{d}\bar{\sigma}^{(0)}} \right) + \mathbf{\tilde{g}}_s^2 \right]_{\mathbb{R}^{1.5}}^{\circ}$ ratio 1.0 Introduce a parametrisation of unknown coefficients in terms of \*ZZ "Theory nuisance parameters"  $\theta$  $\frac{\mathrm{d}\bar{\sigma}_{\mathrm{TNP}}^{(N+1)}}{\mathrm{d}\bar{\sigma}^{(0)}} = \sum^{N} f_{k}^{(j)} \left(\vec{\theta}, x\right) \left(\frac{\mathrm{d}\bar{\sigma}^{(j)}}{\mathrm{d}\bar{\sigma}^{(0)}}\right)$ ddj=1E.g., Bernstein or Chebyshev polynomials







## Fixed Orders & Logs

**Becher:** In processes involving **disparate scales**  $Q_i \gg Q_i$ , higher-order corrections are enhanced by large logarithms  $\alpha_s^n \ln^m(Q_i/Q_i)$ , which can spoil fixed-order truncation. (Max log power  $m \leq 2n$  depends on problem.)









## Shower Logs

How to build an NNLL accurate parton shower

#### ► NLO matching

- ► Correct  $\alpha_s^2$  rate of **neighbour**
- ► Correct  $\alpha_s^2$  rate for a single so  $\alpha_s^2$
- ► Correct  $\alpha_s^2$  rate for a single **c**  $\alpha_s^2$
- ► Correct  $\alpha_s^3$  rate for a single soft-collinear emission



(+ Pythia 8 Hadronization)

 $\mathcal{O}(\alpha_s^3 L^2)$ 

 $\alpha_s^2$ 

 $\alpha_s^2$ 

 $\alpha_s^2$ 

 $e^- \rightarrow Z \rightarrow$  hadrons

 $\sqrt{s} = M_Z = 91.2 \text{ GeV}$ 

 $\alpha_s(M_Z) = 0.118$ 

2-jet@NLO

0.1

0.01

0.

#### Ferrario Ravasio







## Gaps between jets

transverse energy  $E_T$  in gap below  $Q_0$ 





### **Factorization Restoration**

**Factorization** expresses separation of scales

corrections arise – the super-leading logarithms One might think that these effects are numeric hard scale  $Q \sim \sqrt{er}y \frac{\sin all}{hec}$  ause they only arise in higher or but I argue that they can naturally be of the same der as a ode-lbeel coarie bione visitien imperativ study these seffect side add-the correspondence and add-the correspondence add-the correspo corrections to ARisting inference and the calculations. SLLs are caused by a subtle quantum effect: the jet-veto scale @hange of two Coulomb gluons (or Glauber gluons) tween the two initial-state partons in the scatte process, sendiglagatithatibisdeadataion breakdow color coherence, the fact the sum of soft-gluon e sion off two collinear partons has the same effect Soft-collinear factorization violation **Collinear factorization violation** Х by Glauber gluons at  $\mu \sim Q_0$ at  $\mu \sim Q$ a single soft emission off the parent parton. C coherence, however, is the basis for proofs of Q hadronic scale  $\Lambda_{\text{factorization theorems, which underly the theorem$ PDF factorization restored for  $\mu < Q_0$ calculation of all the province sections entry province gives rise to the slept the feat but a breaking rc/·

**PDF Factorization:** long-distance physics contained in universal PDFs ➤ used for all LHC processes — but only proved for (inclusive) Drell-Yan [CSS] Relies on **collinear factorization**. Valid for timelike splittings. Broken for spacelike ones. How bad is it? **Factorization restoration through Glauber gluons:** 



next-to-next-to-leading order (NNNLO) of perturnation observables such as jet cross sections, in which from the jets, the state-of-the-art is NNLO, see

Proof that factorization works at 3-loop order





## $\alpha_{s}$ from hard processes

**Pires:** NNLO  $\alpha_s$  from **dijets** at LHC and HERA Sensitivity from cross section  $\propto \alpha_s^2 \otimes \text{LO} \& \text{PDFs}$ Used: NNLOJET w reduced scale dependence + SLC contributions for the first time for LHC  $\alpha_s$ Central scales:  $\mu_{\text{LHC}}^2 = m_{jj}^2 \& \mu_{\text{HERA}}^2 = Q^2 + \langle p_T \rangle_{1,2}^2$ 

Main Result: "LHC dijets":

 $\alpha_{\rm s}(m_{\rm Z}) = 0.1178 \ (14)_{\rm (fit, PDF)} \ (1)_{(\mu_0)} \ (17)_{(\mu_{\rm R}, \mu_{\rm F})}$ 

LHC + HERA  $\alpha_{\rm s}(m_{\rm Z}) = 0.1180 \,(10)_{\rm (fit, PDF)} \,(1)_{(\mu_0)} \,(22)_{(\mu_{\rm R}, \mu_{\rm F})}$ 

+ Test of RGE running for 7 GeV <  $\mu$  < 7 TeV











## $\alpha_s$ from B and D decays

<b>Che:</b> $\alpha_{\rm s}$ from <b>inclusive semileptonic B decays</b>					
HQE: • $\Gamma\left(B \to X_c \ell \bar{\nu}\right)$ $\Gamma\left(B \to X_c \ell \bar{\nu}_\ell\right)$ =	$\mathcal{E} = \Gamma_0 \begin{bmatrix} C_0 \\ C_0 \end{bmatrix} C$	$C_{\mu_{\pi}} = \frac{C_{\mu_{\pi}}}{C_{\mu_{\pi}}} + \frac{\mu_{\pi}}{2m_{b}^{2}} + \frac{\mu_{\pi}}{2m_{b}^{2}}$	$\frac{1}{2} + C_{\mu_{G}} \frac{\mu_{G}^{2}}{2m_{G}^{2}}$		
$C_0 = \mathbf{c}_0 + \mathbf{c}_1 \frac{\alpha_s}{\pi}$	$+\mathbf{c}_2\left(\frac{\alpha_s}{\pi}\right)$	$)^2 + \mathbf{c}_3 \left( -\frac{\mathbf{c}_3}{2} \right)^2$	$\left(\frac{x_s}{\pi}\right)^3 +$		
(+ indirect sensit	2π. <sup>3</sup> .īvity (?) froi	m running	quark r		
Get IV <sub>cb</sub> I, Guark masses	s from othe	er measure	emensts		
$\Gamma_{sl}$	prediction $[\%]$ o	,(FGeV) [[6]	$\alpha_s \rangle^3$		
$ V_{cb}  = 0.0410 \pm 0.0007$ $\overline{m}_b(\overline{m}_b) = 4.18^{+0.03}_{-0.02} \text{ GeV}$ $\overline{m}_b(\overline{m}_b) = 1.27 \pm 0.02 \text{ CeV}$	3.4 (1.4) 3.0 (1.1) 2.1 (1.4)	3.1 (1.3) 2.7 (1.0) 1.8 (1.2)	$\left(\frac{1}{\pi}\right)$		
$m_c(m_c) = 1.27 \pm 0.02 \text{ GeV}$ R-scale $\mu = 5^{+5}_{-2.5} \text{ GeV}$ High order power corrections	4.4(2.2)	4.0(2.0)	Cor		
$\begin{aligned} \tau_{B^{\pm}} &= 1.638 \pm 0.004  \mathrm{ps} \\ \mathcal{B}(B^{\pm} \to X_c \ell \nu) &= 10.8 \pm 0.4  \% \end{aligned}$	-	$0.0 \\ 0.2 \\ 2.4 (1.8)$	$\Delta \alpha_s M_s$		
Sum	6.7(3.2)	6.5(3.4)	pc		
	$\overline{m}_b (5 \text{ Ge})$	$\overline{W} \overline{m}_{c}(5)$	GeV)		
riond QCD 202					

### B

#### Che, Wu



$$\vdash \mathcal{O}(\alpha_s^4)$$

$$petitive = 0.0018$$

$$\alpha_s(5 \text{ GeV})$$

$M_{i}: \alpha_{s}$ from inclusive semileptonic D dec $\alpha_{s}(5 \text{ GeV})$ Understanding of a ct of spectator quark in 0.17, 0.18, 0.19,         semileptonic D decays $\Gamma_{SL}$ $\Gamma_{SL}$ $\sigma_{s}(5 \text{ GeV})$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(5 \text{ GeV})$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(5 \text{ GeV})$ $D_{l}$ $\mathcal{B}_{SL}$ (%) $\tau$ (10 <sup>-13</sup> s) $\Gamma_{SL}$ (10 <sup>-15</sup> G $D^{0}$ $6.46 \pm 0.09 \pm 0.11$ $4.10 \pm 0.01$ $104 \pm 2$ $D^{+}$ $6.30 \pm 0.13 \pm 0.10$ $5.04 \pm 0.04$				
$\begin{array}{c} u_{s}(5 \text{ GeV}) \\ \text{Jn} \underbrace{ u_{s}(5 \text{ GeV})}_{0.17, \ 0.18, \ 0.19, \ldots} \\ \text{Semileptonic D decays} \\ \hline \Gamma_{SL} \\ \text{Different BRs and lifetimes:} \\ \hline \overline{m_{L}}(\overline{m_{L}}) \\ \hline D_{i} \\ \mathcal{B}_{SL}(\%) \\ \hline D^{0} \\ 0 \\ 6.46 \pm 0.09 \pm 0.11 \\ 104 \pm 0.01 \\ 104 \pm 2 \\ D^{+} \\ 16.13 \pm 0.10 \pm 0.29 \\ \hline D_{s} \\ 6.30 \pm 0.13 \pm 0.10 \\ \hline 5.04 \pm 0.04 \\ \hline 82 \pm 2 \\ \hline \end{array}$	<b>Wμ</b> : α	, trom inclusive	semilepto	nic D dec
$\begin{array}{c c} \text{Jn} \begin{array}{c} \begin{array}{c} \begin{array}{c} \begin{array}{c} \text{Diff} \\ \textbf{erent BRs and lifetimes:} \\ \hline m_{L}(\overline{m}_{L}) \end{array} \end{array} \\ \begin{array}{c} \begin{array}{c} \begin{array}{c} \\ \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \begin{array}{c} \begin{array}{c} \\ \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \begin{array}{c} \begin{array}{c} \\ \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \begin{array}{c} \begin{array}{c} \\ \end{array} \\ \hline \end{array} \\ \begin{array}{c} \begin{array}{c} \\ \\ \end{array} \\ \hline \end{array} \\ \begin{array}{c} \begin{array}{c} \\ \\ \end{array} \\ \hline \end{array} \\ \begin{array}{c} \begin{array}{c} \\ \\ \end{array} \\ \hline \end{array} \\ \\ \hline \end{array} \\ \\ \hline \end{array} $ \\ \hline \end{array}  \\ \hline \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \hline \end{array} \\ \\ \end{array} \\ \hline \end{array} \\ \hline \end{array}  \\ \hline \end{array}  \\ \hline \end{array} \\ \end{array} \\ \hline \end{array} \\ \\ \end{array}  \\ \hline \end{array} \\ \hline \end{array}  \\ \hline \end{array} \\ \end{array}  \\ \hline \end{array}  \\ \hline \end{array} \\ \end{array} \\ \end{array} \\ \hline \end{array} \\ \hline \end{array}  \\ \hline \end{array}  \\ \hline \end{array} \end{array} \\ \end{array}  \\ \hline \end{array} \end{array} \\ \end{array} \\ \end{array}  \\ \hline \end{array} \\ \end{array} \end{array} \\ \end{array} \end{array} \\ \end{array}  \\ \end{array}  \\ \hline \end{array}  \\  \\		$a_s(J \cup e v)$	$\langle \rangle$	
$\frac{1}{m_{L}(m_{L})} = \frac{1}{m_{L}(m_{L})} = \frac{1}{m_{L}(5 \text{ GeV})}$ Semileptonic D decays $\frac{1}{SL} = \frac{1}{m_{L}(m_{L})} = \frac{1}{m_{L}(5 \text{ GeV})}$ $\frac{1}{m_{L}(m_{L})} = \frac{1}{m_{L}(5 \text{ GeV})} = \frac{1}{m_{L}(5 \text{ GeV})}$ $\frac{1}{m_{L}(m_{L})} = \frac{1}{m_{L}(m_{L})} = \frac{1}{m_{L}(10^{-13} \text{ s})} = \frac{1}{m_{L}(10^{-15} \text{ GeV})}$ $\frac{1}{m_{L}(10^{-15} \text{ GeV})} = \frac{1}{m_{L}(10^{-15} \text{ GeV})}$	Under	rstandingnosact of	f spetchator	quark in
Different BRs and lifetimes: $c, u$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(5 \text{ GeV})$ $D_{i}$ $\mathcal{B}_{SL}$ (%) $\tau$ (10 <sup>-13</sup> s) $\Gamma_{SL}$ (10 <sup>-15</sup> G $D^{0}$ $6.46 \pm 0.09 \pm 0.11$ $4.10 \pm 0.01$ $104 \pm 2$ $D^{+}$ $16.13 \pm 0.10 \pm 0.29$ $10.33 \pm 0.05$ $103 \pm 2$ $D_{s}^{+}$ $6.30 \pm 0.13 \pm 0.10$ $5.04 \pm 0.04$ $82 \pm 2$		$\frac{0.17, 0.18, 0.19,}{0.100000000000000000000000000000000$		•
Different BRs and lifetimes: $c, u$ $\overline{m}_{L}(\overline{m}_{L})$ $\overline{m}_{L}(5 \text{ GeV})$ $D_{i}$ $\mathcal{B}_{SL}$ (%) $\tau$ (10 <sup>-13</sup> s) $\Gamma_{SL}$ (10 <sup>-15</sup> G $D^{0}$ $6.46 \pm 0.09 \pm 0.11$ $4.10 \pm 0.01$ $104 \pm 2$ $D^{+}$ $16.13 \pm 0.10 \pm 0.29$ $10.33 \pm 0.05$ $103 \pm 2$ $D_{s}^{+}$ $6.30 \pm 0.13 \pm 0.10$ $5.04 \pm 0.04$ $82 \pm 2$	Sennie	$\Gamma_{a}$		
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$D^+$ 16.13 $\pm 0.10 \pm 0.29$ 10.33 $\pm 0.05$ 103 $\pm 2$ $D_s^+$ 6.30 $\pm 0.13 \pm 0.10$ 5.04 $\pm 0.04$ 82 $\pm 2$	$D^0$	$6.46 \pm 0.09 \pm 0.11$	$4.10\pm0.01$	$104 \pm 2$
$D_s^+$ 6.30 ± 0.13 ± 0.10 5.04 ± 0.04 82 ± 2				102 + 2
	$D^+$	$16.13 \pm 0.10 \pm 0.29$	$10.33 \pm 0.05$	$103 \pm 2$
$\boldsymbol{\nu}$	$D^+$ $D^+_{\epsilon}$	$16.13 \pm 0.10 \pm 0.29$ $6.30 \pm 0.13 \pm 0.10$	$10.33 \pm 0.05$ 5 04 + 0 04	$103 \pm 2$ 82 + 2

Same  $\Gamma_{SL} \xrightarrow{f}{s}$  spectator<sup>2</sup> impact is  $G^{\pm}$  in symmetric in semileptonig decays?

$m_c[{ m GeV}]$	1.3
$lpha_{S}(m_{c}^{2})[10^{-3}]$	44
S I	

$$\begin{array}{c}
 D^+, D^0 \\
 3701 \pm 0.0338 \\
 445 \pm 9 \pm 114 \\
 4
 \end{array}$$



 $\ell^\pm$ 





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 $D_s^+$ 

## $\alpha_{s}$ from lattice

#### **Review of Flavour Lattice Averaging Group (FLAG 2024)**

#### **Overview of** $\alpha_{c}$ **determinations:**

a); OnDeter the affer the side of stealed and the Berth the Reasonable resolution for big box Cannot go higher than this  $\mu a \ll 1, L = N_s a > 1 \text{ fm} \Rightarrow \mu = 1 - 3 \text{ GeV}$ • Comparison with lattice perturbation theory  $(\mu = 1/a) \xrightarrow{\sim} \alpha^{\text{MS}} \alpha^{\text{MS}} (\mu^{1/a}) (\mu^{1/a}) (a)$ = 1/aLimited by accuracy of lattice perturbation theory

• Lattice OCD in "femto boxes" and special schemes  $\alpha_s^{SF}(\mu = 1 L)$ lculations Katel and the step scaling to the stand of the store of the store of the step scaling step scaling the step scaling and the step scaling step scale and the step

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#### Theory Summary — Peter Skands

## PDFs from lattice

The natural ab-initio method to study QCD non-perturbatively is on the **lattice**. But ...

PDFs  $\equiv$  expectation values of bilocal operators evaluated along **light-like** lines.

Cannot evaluate this on a Euclidean setup
→ Traditional lattice studies were limited to first few (three) Mellin moments of a local ME.

**Breakthrough** (Ji 2013): put quarks some distance apart. Then boost them to heck → almost lightlike separation (in proton frame).

→ Use perturbative matching from finite to infinite momentum (and deal with divergencies)

#### Zafeiropoulos



Take-home: Lattice can by now provide ab-initio PDF determinations without theoretical obstructions



## Inclusive $b \rightarrow s \mu^+ \mu^-$

• Exclusive  $b \rightarrow s\mu\mu$  subject to potentially large and uncontrolled power corrections



• Up to power corrections the inclusive rate is free of hadronic uncertainties:

$$\Gamma[B \to X_s \ell \ell] = \Gamma[b \to X_s \ell \ell] + O\left(\frac{\Lambda_{QCD}^2}{m_b^2}, \cdots\right)$$

but not for the ratio:

 $\mathscr{R}(q_0^2) =$ 

- $q^2 > 14.4 \text{ GeV}^2$
- prediction for  $\mathscr{R}$ :

 $\mathscr{B}[>15]_{\text{SM+Belle}} = (4.10 \pm 0.81) \times 10^{-7}$ 

### Lunghi

• The OPE breaks down at large  $q^2 = m_{\mu\mu}^2$ 

$$= \frac{\int_{q_0^2}^{m_b^2} dq^2 \frac{\mathrm{d}\Gamma(\bar{B} \to X_s \ell^+ \ell^-)}{dq^2}}{\int_{q_0^2}^{m_b^2} dq^2 \frac{\mathrm{d}\Gamma(\bar{B}^0 \to X_u \ell \nu)}{dq^2}}$$

• LHCb has already measured enough high- $q^2$  modes to reconstruct the BR:  $\mathscr{B}[>15]_{||HCb|} = (2.65 \pm 0.17) \times 10^{-7}$ 

• LHCb should produce a proper combination of these modes taking into account correlations and for

• Using differential Belle semileptonic data the high- $q^2$  and the theory

- Inclusive modes are currently in agreement with data
- Future LHCb + Belle II data can confirm the exclusive anomalies at  $5\sigma$  if central values do not change:









## Hadronic Transitions & Form Factors

#### Mishra: $B \rightarrow K$ FFs from LCSR

Measured BR( $B \rightarrow K \mu \mu$ ) lower than predictions. Tensions between different theory predictions.



Can safely neglect non-fact. soft-gluon contributions to c loop

+ Re-analysis of hadronic ME -> tension persists

d'Ambrosio:  $K \rightarrow \pi \ell^+ \ell^-$ Dominated by long distances **LC**  $\Rightarrow$  Sum of 1-meson poles  $\Rightarrow$  VMD-like ansatz for  $\gamma^*$  $\langle J(q)J(-q)\rangle = \sum \frac{a_n^2}{q^2 - m_n^2} \widetilde{q_2 \to \infty} \log q^2$ 

-

#### **Gubernari:** improving FF parametrisations

Conventional approach: BGL

⇒ divergent series in presence of branch cuts

Problem for  $B \to K, B \to D^{(*)}, \Lambda_h \to \Lambda$ 

(Also: truncation error meaningless)

Analytic structure suggests an alternative parm:



(+ extension to rescattering)





## FFs for B<sub>c</sub> & HF hadrons from recombinations?

### Nandi: FFs for $B_c \rightarrow$ charmonium

Tests of Lepton Flavour Universality analogous to those in  $B \rightarrow D^{(*)}$  (R<sub>D</sub>, R<sub>D\*</sub>) can be done with  $B_c \to \eta_c, B_c \to J/\psi$ , and  $B_c \to P(=\chi_c^0, \chi_c^1, h_c)$ For  $B_c \rightarrow J/\psi$ , there are FFs from lattice Heavy Quark Spin Symmetry  $\Rightarrow B_c \rightarrow \eta_c$ Estimate symmetry-breaking correction ~ 30%

 $\Rightarrow R(\eta_c) = 0.290 \pm 0.017$ 

Also compute  $B_c \rightarrow P$ - and S-wave FFs in both NRQCD ( $q^2 \rightarrow 0$ ) & pQCD (high  $q^2$ )

In: B<sub>c</sub>-> J/ $\psi$  FFs, decay consts, and measured charmonium radiative decay rates  $\Rightarrow$  LFU R(...)

### Li:

- **Observation**: more low-z D\* in data than in baseline MCs
- Could originate from **recombination** of c with quark from UE?
- Similar to coalescence?



- Assumptions for "UE sea"  $\rightarrow$  good fits to data
- **Also:** more HF baryons at low  $p_T$  than at LEP

Diquark-style recombinations?

(Note: similar phenomena modelled in event generators: **QCD** Colour Reconnections)

→ can also make predictions for tetra-quarks







## Pion Holography?

**Holography** = formal equivalence beween two theories, said to be each other's **holographic duals**:

- A strongly-coupled scale-invariant (conformal) gauge theory in flat 4D space-time
- A weakly-coupled 5D gravity (string) theory in a curved space (AdS)
- ⇒ can do calculations in the weakly coupled theory and relate them to the stronglycoupled one!

#### **Problems:**

- QCD  $\neq$  scale invariant
- Longitudinal dofs neglected in "light-front quantisation"  $\Rightarrow$  massless pion



**Sandapen:** restoring conformal-symmetry breaking longitudinal potential  $U_{\parallel} \rightarrow \text{correct } m_{\pi}$ 

**Three Different** forms of  $U_{\parallel} \Rightarrow$  same  $m_{\pi'} f_{\pi'} r_{\pi'}$ , low-Q2 form factors, and  $\Gamma_{\gamma\gamma}$  = 7.0, 7.2, and 7.4 eV

PDG:  $7.82 \pm 0.22 \text{ eV}$ 

Form "C" with  $\Gamma_{\gamma\gamma} = 7.4 \text{ eV}$ also exhibits quantitative agreement with holographic prediction in limit of weak coupling

#### Work in progress:

holographic pion distribution amplitude & pion PDFs























Kolbé: Simulating jets in medium

Nonzero  $v_2$  at high  $p_T$  not well understood



### Largest TH uncertainty: hadronic vacuum polarisation (HVP)

- 1. Data driven (from R-ratio)
  - $e^+e^- \rightarrow had + optical theorem$
  - But note some tensions among data sets!
- 2. Lattice QCD (10<sup>10</sup>-dim integral)
  - Community agreement on intermediate (simpler) benchmark: **window observable** (restrict correlator to 0.4-1fm)
  - Finer lattice spacings (page 10)

#### 3. Frankenfit

- Lattice for 0-2.8fm window
- **Data driven** for  $2.8-\infty$  fm tail (5% of total result, avoids  $\rho$  peak, good agreement).

#### Tóth





#### Apologies: a few **speakers** not yet mentioned in these slides

### **Quark Flavour Physics**

- S. Nandi ( $B_c \rightarrow X_{c\bar{c}}$  FFs)
- E. Lunghi  $(b \rightarrow s\mu^+\mu^-)$
- N. Gubernari ( $B \rightarrow K FFs$ )
- D. Mishra ( $B \rightarrow K LCSR$ )
- G. d'Ambrosio (LC  $K \rightarrow \pi \ell^+ \ell^-$ )

### **Non-perturbative QCD**

- S. Li (Hadrons from recombination)
- S. Zafeiropoulos (Lattice PDFs)
- B. Tóth (Lattice  $g_{\mu} 2$ )
- M. Praszalowicz (Elastic pp)
- R. Sandapen (Pion Holography)

### **Heavy lons**

- N. Zardishti (HI Overview)
- I. Kolbé (Jets in Medium)
- E. Speranza (Spins in HI)
- W. Schee (QCD ↔ Gravity?)

### Methodology

- R. Poncelet ( $\sigma_{\rm th}$ )
- M. White (Q. Advantage in  $t\bar{t}$ )



EW/BSM  $Q \gtrsim M_{\rm EW}$ Ŋ Fixed-Order  $\checkmark$ pQCD Physics Single-Scale Matching  $Q \gg \Lambda$ Flavour Resummed pQCD Multi-Scale  $Q_i \gg Q_j \gg \Lambda$ Hadronization Strings, QGP  $Q \sim [\Lambda, 5\Lambda]$ Hadrons  $Q \lesssim \Lambda$  $g_{\mu} - 2$ 

### Hard Processes

- S. Jaskiewicz ( $m_t$  in HH)
- M. Kerner  $(m_t \text{ in } H + j)$
- M. Marcoli (pp  $\rightarrow \gamma \gamma$ )
- D. d'Enterria (Rare *H* decays)

### **Splittings, Resummation, Factorization**

- A. Pelloni (N3LO splittings)
- F. Hekhorn (aN3LO PDFs & Higgs)
- S. Ferrario Ravasio (Shower Logs)
- N. Schalch (Non-global Logs)
- T. Becher (Super-Leading Logs)
- M. Neubert (Factorization)
- A. Vladimirov (DY TMD)
- W. Zhan (DY intrinsic  $p_T$ )

### **BSM**

- S. Vempati (Higgs FLV)
- E. Vryonidou (Higgs CPV)
- M. Baker (Heavy Vectors)
- S. Balan (Dark-Matter Fits)
- M. Fedele (Sterile  $\nu$  in  $b \rightarrow c\ell^+\ell^-$ )

### **Strong Coupling**

- J. Pires ( $\alpha_s$  from Dijets)
- M. Benitez ( $\alpha_s$  from M<sub>H</sub>)
- Y. Che ( $\alpha_s$  from B)
- J. Wu ( $\alpha_{\rm s}$  from D)
- P. Petreczky ( $\alpha_s$  from lattice)





#### Comments (04-Apr-2025 18:29:14) Machine checkout until Monday

Transfer Line tests completed

#### AFS: Single\_10b\_4\_2\_4

### Theory Summary 2025

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59<sup>th</sup> Rencontres de Moriond: QCD & High Energy Interactions, La Thuile, April 2025

