# **QFT with Hadrons**Introduction to B Physics

→ 1. Leptonic Decays of Hadrons: from  $\pi$  →  $\ell \nu$  to B →  $\ell \nu$ 

QFT in Hadron Decays. Decay Constants. Helicity Suppression in the SM.

2. On the Structure and Unitarity of the CKM Matrix

The CKM Matrix. The GIM Mechanism. The Unitarity Triangle.

3. Semi-Leptonic Decays and the "Flavour Anomalies"

 $B \to D^{(*)} \ell v$ . The Spectator Model. Form Factors. Heavy Quark Symmetry.

 $B \to K^{(*)} \ell^+ \ell^-$ . FCNC. Aspects beyond tree level. Penguins. The OPE.

# Recap of (applied) QFT

#### Want to:

Start from assumed field content & Lagrangian (e.g., SM).

Compute scattering cross sections and decay rates.

Total and differential

Compare to experimental measurements.

#### Recipe in perturbative QFT:

Set up (relativistically normalised) in- and outgoing states.

Interaction picture: plane-wave states (eigenstates of free theory, in momentum space)

Compute (Lorentz-invariant) transition amplitudes.

QFT under the hood: Dyson's Formula, Wick Contractions

→ For practical calculations: **Feynman rules & diagrams** 

Sum over amplitudes, square, and keep terms to given perturbative order.

Integrate over the relevant (Lorentz-invariant) **phase space**(s).

#### Recap: Decay Rates

Partial decay rate (a.k.a., "partial width") of particle of mass M into n bodies, in its CM:

$$\Gamma_{i \to f}$$

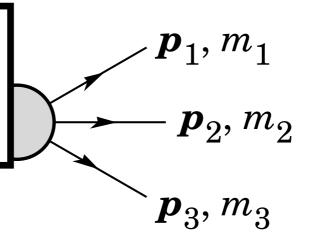
**Total Width** = sum over partial widths

$$\Gamma_i = \sum_j \Gamma_{i \to j}$$

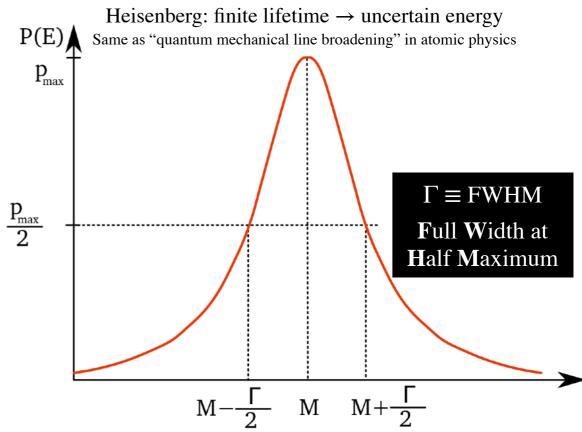
**Average Lifetime** 

$$\tau = 1/\Gamma$$

=  $\hbar/\Gamma$  if not using natural units

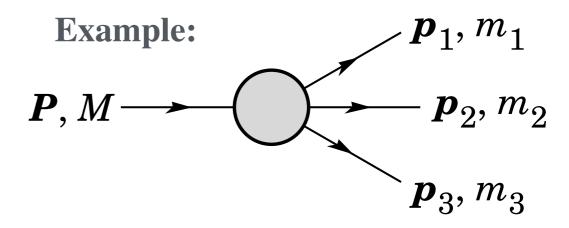


#### Why is it called the width?



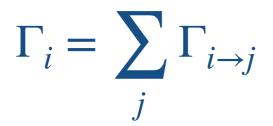
### Recap: Decay Rates

Partial decay rate (a.k.a., "partial width") of particle of mass M into n bodies, in its CM:





Total Width = sum over partial widths



#### **Average Lifetime**

$$\tau = 1/\Gamma$$

=  $\hbar/\Gamma$  if not using natural units

#### Branching fractions = $\Gamma_j/\Gamma$

Example:  $\pi^+$  decays (see, e.g., pdg.lbl.gov)

$$BR(\pi^+ \to \mu^+ \nu_\mu)$$
 (99.98770±0.00004) %   
  $BR(\pi^+ \to e^+ \nu_e)$  (1.230 ±0.004) ) × 10<sup>-4</sup>

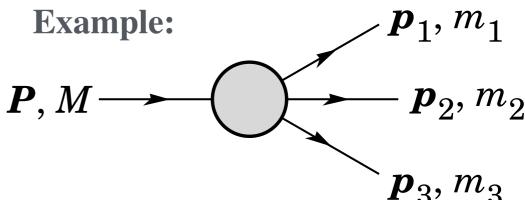
This agrees with the SM prediction.

Our first application: weak leptonic decays of hadrons

### Recap: Decay Rates

#### **Reminder:**

Fermi's Golden Rule (on relativistic form) for decay rates:



$$\Gamma_{i\to f} = \int d\Gamma_{i\to f} = \frac{(2\pi)^4}{2M} \int |\mathcal{M}_{i\to f}|^2 d\Phi_n(P; p_1, ..., p_n)$$

Lorentz-invariant Matrix Element

Lorentz-invariant **phase-space element**:

$$d\Phi_n(P; p_1, \dots, p_n) = \delta^4 \left(P - \sum_{i=1}^n p_i\right) \prod_{i=1}^n \frac{d^3 p_i}{(2\pi)^3 2E_i}$$

**a.k.a.** : *dLIPS* 

=  $d^4p_i$  with on-shell condition (**L.I.**)

(see, e.g., PDG review (pdg.lbl.gov) section 47: kinematics)

# Special case: 2-body decays

In 2-body decays, the kinematics are fully constrained (up to an overall solid angle)



$$\Rightarrow \Gamma_{i \to f} = \frac{|\mathbf{p}^*|}{32\pi^2 M^2} \int |\mathcal{M}_{fi}|^2 \mathrm{d}\Omega \qquad \begin{array}{c} \text{VALID FOR } \underline{\text{ALL}} \\ \text{2-BODY DECAYS} \\ \text{Exercise problem E1a: derive this} \end{array}$$

**Exercise problem E1a:** derive this formula from the one on the previous page.

with  $\mathbf{p}^*$  the 3-momentum of either of the decay products in the rest frame of M:

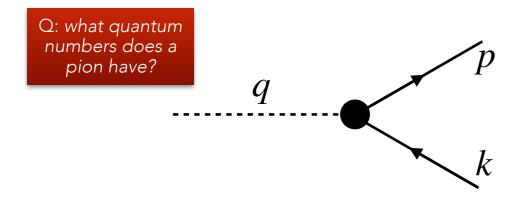
$$p^* = \frac{1}{2M} \sqrt{[M^2 - (m_1 + m_2)^2][M^2 - (m_1 - m_2)^2]}$$

Question: why does it not matter which 3-momentum we use?

**Exercise problem E1b:** derive this formula for p\*

# OK, let's apply this to compute pion decays

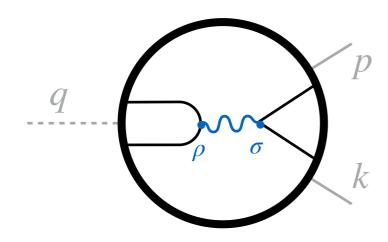
#### Want to calculate M for: $\pi^-(q) \to \mu^-(p) + \bar{\nu}_{\mu}(k)$



First problem: the SM Lagrangian does not include a "pion"

How are we supposed to apply Feynman rules without a  $\pi$ - $\mu$ - $\nu$  vertex?

#### What is really going on?



It's the **weak force**: W exchange between quark and lepton currents

$$m_{\pi} = 0.13 \text{ GeV}$$
  
 $q = (m_{\pi}, 0, 0, 0)$   
 $m_{W} = 80.4 \text{ GeV}$ 

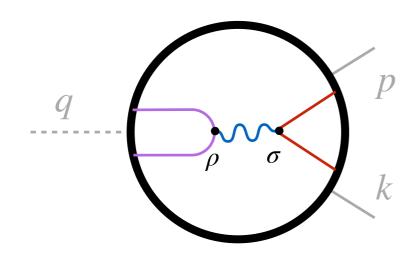
W propagator: (how) familiar is this?

$$\frac{-i(g_{\rho\sigma} - q_{\rho}q_{\sigma}/M_W^2)}{q^2 - M_W^2} \rightarrow \frac{ig_{\rho\sigma}}{M_W^2}$$

# Application to Pion Decay

# Want to calculate M for: $\pi^-(q) \to \mu^-(p) + \bar{\nu}_{\mu}(k)$

What is really going on?



$$m_{\pi} = 0.13 \text{ GeV}$$
  
 $q = (m_{\pi}, 0, 0, 0)$   
 $m_{W} = 80.4 \text{ GeV}$ 

W propagator: 
$$\frac{ig_{\rho\sigma}}{M_W^2}$$

**Lepton current:** 
$$L^{\sigma}(p,k) = -i \frac{g_w}{2\sqrt{2}} \bar{u}(p) \gamma^{\sigma} (1-\gamma_5) v(k)$$
 (how) familiar is this?

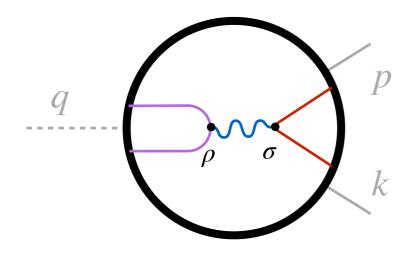
Quark current: 
$$\frac{-i\frac{g_w}{2\sqrt{2}}\bar{v}_u\gamma^\rho(1-\gamma_5)u_d}{2\sqrt{2}}$$
 Why not?

Monash

#### The Quark Current

#### The quark-antiquark pair

Bouncing around inside the pion  $\rightarrow$  not free plane-wave states.



$$\mathcal{M}(\pi o \mu \bar{
u}) = Q^{
ho}(q) \frac{ig_{
ho\sigma}}{M_W^2} L^{\sigma}(p,k)$$

#### What do we know about the quark current?

Must be proportional to  $g_w$ 

Carries a 4-vector index, o

Since the pion has spin 0 (no spin vector), the only 4-vector is: q

$$\Rightarrow Q^{\rho}(q) = \frac{g_w}{2\sqrt{2}} q^{\rho} f(q^2)$$

$$\int q^2 = m_{\pi}^2 = \text{const.}$$

$$= \frac{g_w}{2\sqrt{2}} q^{\rho} f_{\pi}$$
f . "Pion decay constant"

University

# So the matrix element for $\pi^-(q) \to \mu^-(p) + \bar{\nu}_{\mu}(k)$ is:

$$G_{F} = \frac{\sqrt{2}g_{w}^{2}}{8M_{W}^{2}} \qquad \mathcal{M} = \frac{G_{F}}{\sqrt{2}} (p^{\rho} + k^{\rho}) f_{\pi} \left[ \bar{u}(p) \gamma_{\rho} (1 - \gamma_{5}) v(k) \right]$$

Use the Dirac eqs. for the <u>neutrino</u> and <u>muon</u>:

$$kv(k) = 0 \qquad \bar{u}(p)(\not p - m_{\mu}) = 0$$

Only a term proportional to the muon mass survives

$$\mathcal{M} = \frac{G_F}{\sqrt{2}} f_{\pi} m_{\mu} \bar{u}(p) (1 - \gamma_5) v(k)$$

$$\implies |\mathcal{M}|^2 = \frac{G_F^2}{2} f_\pi^2 m_\mu^2 \operatorname{Tr} \left[ (\not p + m_\mu) (1 - \gamma_5) \not k (1 + \gamma_5) \right]$$

$$=8(p \cdot k)$$
 (how) familiar is this?

## Putting it Together

From previous slide:

$$|\mathcal{M}|^2 = 4G_F^2 f_\pi^2 m_\mu^2 (p \cdot k)$$

We also had the Golden-rule master formula for  $1\rightarrow 2$  decays

$$\Gamma_{i\to f} = \frac{|\mathbf{p}^*|}{32\pi^2 M^2} \int |\mathcal{M}_{fi}|^2 d\Omega$$

with 
$$p^* = \frac{m_\pi}{2} \left( 1 - \frac{m_\mu^2}{m_\pi^2} \right)$$
 cf. your derivation of p\*

 $\rightarrow p = (E_{\mu}, p^*)$ and  $k = (p^*, -p^*)$   $q = (m_{\pi}, 0, 0, 0)$   $\implies (k \cdot p) = (k \cdot (q - k))$  $= m_{\pi} |p^*|$ 

## $\Gamma(\pi \rightarrow \mu \nu)$

$$\Rightarrow \Gamma(\pi \to \mu \bar{\nu}) = \frac{G_F^2}{8\pi} f_\pi^2 m_\pi m_\mu^2 \left( 1 - \frac{m_\mu^2}{m_\pi^2} \right)^2$$

Question: could we use same  $G_F$  for  $\Gamma(\pi \rightarrow ev)$ ? Same  $f_{\pi}$ ?

Can get  $G_F$  from muon decay (no hadrons  $\triangleright$  no decay constant).

But cannot compute  $f_{\pi}$  (perturbatively), so cannot "predict" pion lifetime.

Instead, we can use the pion lifetime to **measure**  $f_{\pi}$ .

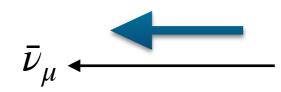
 $m(\pi,\mu,e,v) = (135, 105, 0.5, 0) MeV$ 

Independently of  $f_{\pi}$  however, we can now account for:

$$BR(\pi^+ \to \mu^+ \nu_\mu)$$
  
 $BR(\pi^+ \to e^+ \nu_e)$ 

$$(99.98770\pm0.00004)\%$$
  
 $(1.230\pm0.004)\times10^{-4}$ 

Physics = Angular momentum cons.:



Spin 0



(how) familiar is this?

In SM,  $\bar{\nu}$  is massless and **right-handed**  $\Rightarrow$  **positive helicity** 

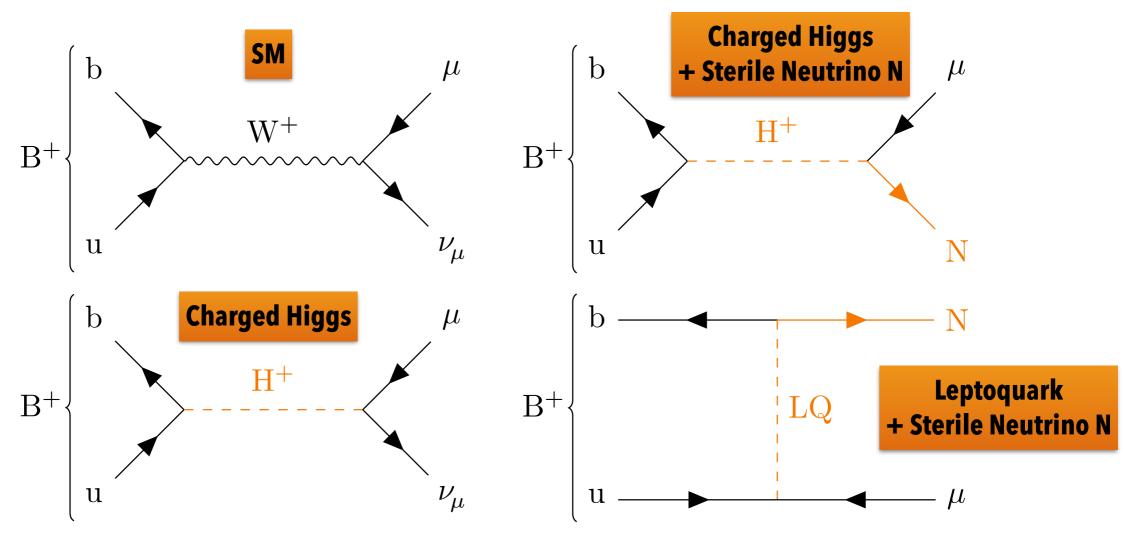
⇒ Muon must also have positive **helicity**, but W couples to left-handed **chirality**.

 $\langle u_L | u_+ \rangle \propto m \Leftrightarrow \text{Helicity Suppression}$ 

## $B^+ \rightarrow \tau^+ v$ and $B^+ \rightarrow \mu^+ v$

#### A very similar treatment applies to B+ $\rightarrow$ $\tau$ + $\nu$ and B+ $\rightarrow$ $\mu$ + $\nu$

Some reasons why those might be interesting: (illustration from arXiv:1911.03186)



Most BSM diagrams not helicity suppressed! (why?)  $\Rightarrow$  Can be even larger than SM amplitude despite heavier virtual states. (BSM currents not restricted to be purely L-handed)

Exercise problem E3: give reason(s) why B decays might be more interesting than pion decays?

## Research Problems for Assignment

# R1. Provide an elaborate derivation of $\mathcal{M} \Leftrightarrow |\mathcal{M}|^2 \Leftrightarrow \Gamma$ $\Rightarrow$ Branching Fraction for B+ $\rightarrow \tau$ + $v_{\tau}$ in the SM and compare with measurements

Use the lattice determination of  $f_B$  from <a href="https://arxiv.org/abs/1607.00299">https://arxiv.org/abs/1607.00299</a>

Use the Heavy-Flavour Averaging Group (HFLAV) value for V<sub>ub</sub> from <a href="https://arxiv.org/abs/1909.12524">https://arxiv.org/abs/1909.12524</a>

Find measured values for the lifetime of the B+ meson and BR(B+ $\rightarrow \tau$ + $\nu_{\tau}$ ) in the Particle Data Group (PDG) summary for the B+ meson: <u>pdg.lbl.gov</u>

(You will also need the masses of the involved particles, and the value of the Fermi constant,  $G_F$ )

#### R2. What is BR(B+ $\rightarrow \mu^+\nu_\mu$ )/BR(B+ $\rightarrow \tau^+\nu_\tau$ ) in the SM?

Belle has reported a measurement of BR(B+ $\rightarrow$  $\mu$ + $\nu_{\mu}$ ), see https://arxiv.org/abs/1911.03186: study it, and does it agree with your expectation?

# Summary of Problems and Exercises for Home Study

- E1. Derive the formulae for  $\Gamma_{1\rightarrow 2}$  & p\* on p.5. You may use standard textbooks such as Thomson / Griffiths / Halzen & Martin / ...
- E2. Perform the detailed steps in the derivation on p.9
- E3. Give reason(s) why B decays may be more interesting than  $\pi$  ones?

You will present your progress on these in the next lesson and we will discuss any questions / issues you encounter.

+ Assignment Problems 1&2: the B physics research problems on p.14

Due in week 6